



classmate



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QUANTUM SEARCH ALGORITHMS

Suppose we are given a map containing many cities, and wish to determine the shortest route passing thro' all cities on the map. A simple algorithm to find this route is to search all possible routes thro' the cities, keeping a running record of which route has the shortest length.

On a classical computer, if there are N possible routes, it takes $O(N)$ operations to determine the shortest route using this method.

There is a quantum search algorithm, sometimes known as Grover's algorithm, which enables this search method to be sped up substantially, requiring only $O(\sqrt{N})$ operations.

□ The Quantum Search Algorithm (Grover's algorithm)

□ The oracle

Suppose we wish to search thro' a search space of N elements. Rather than search the elements directly, we concentrate on the index to those elements, which is just a number in the range 0 to $N-1$.

For convenience we assume $N=2^n$, so the index can be stored in n bits, and that the search problem has exactly M solutions, with $1 \leq M \leq N$.

A particular instance of the search problem can conveniently be represented by a function f , which takes as i/p an integer x , in the range 0 to $N-1$. By definition,

$f(x) = 1$ if x is a solution to the search problem.

$f(x) = 0$ if x is not a solution to the search problem.

Suppose,

we are supplied with a quantum oracle, with the ability to recognize solutions to the search problem. This recognition is signalled by making use of an oracle qubit.

The oracle is a unitary operator $\mathcal{O}$, defined by its action on the computational basis:

$$|x\rangle|q\rangle \xrightarrow{\mathcal{O}} |x\rangle|q \oplus f(x)\rangle$$

where $|x\rangle$ is the index register

$\oplus$ denotes addition modulo 2

and the oracle qubit $|q\rangle$ is a single qubit which is flipped if $f(x)=1$, and is unchanged otherwise.

We can check whether x is a solution to our search problem by preparing $|x\rangle|0\rangle$, applying the oracle, and checking to see if the oracle qubit has been flipped to $|1\rangle$.

In the quantum search algorithm it is useful to apply the oracle with the oracle qubit initially in the state $\frac{|0\rangle - |1\rangle}{\sqrt{2}}$, just as was done in the Deutsch-Jozsa algorithm.

If x is not a solution to the search problem, applying the oracle to the state $|x\rangle \frac{|0\rangle - |1\rangle}{\sqrt{2}}$ does not change the state.

If x is a solution to the search problem, then $|0\rangle$ and $|1\rangle$ are interchanged by the action of the oracle, giving a final state $-|x\rangle \frac{|0\rangle - |1\rangle}{\sqrt{2}}$.

The action of the oracle is:

$$\begin{aligned}
 |x\rangle \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) &\xrightarrow{O} |x\rangle \left(\frac{|0\rangle \oplus |f(x)\rangle - |1\rangle \oplus |f(x)\rangle}{\sqrt{2}} \right) \\
 &= |x\rangle \left(\frac{|f(x)\rangle - |f(x)\rangle}{\sqrt{2}} \right) \\
 &= (-1)^{f(x)} |x\rangle \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right)
 \end{aligned}$$

The state of the oracle qubit is not changed, which remains $\frac{|0\rangle - |1\rangle}{\sqrt{2}}$ throughout the quantum search algorithm, and therefore be omitted from further discussion of the algorithm.

With this convention, the action of the oracle may be written:

$$|x\rangle \xrightarrow{O} (-1)^{f(x)} |x\rangle$$

⇒ The Oracle marks the solutions to the search problem, by shifting the phase of the solution.

There is a distinction b/w knowing the solution to a search problem, and being able to recognize the solution. The crucial point is that it is possible to recognize the solution without necessarily being able to know the solution.

- For an N item search problem with M solutions, we need to only apply the search oracle $O(\sqrt{N/M})$ times in order to obtain a solution, on a quantum computer.

- There are many computational problems in which it's difficult to find a solution, but relatively easy to verify a solution.

Ex:- We can easily verify a solution to a sudoku by checking all the rules are satisfied. For these problems, we can create a function f that takes a proposed solution x , and returns $f(x)=0$ if x is not a solution ($x \neq z$) and $f(x)=1$ for a valid solution ($x=z$). Our oracle can then be described as:

$$O|x\rangle = (-1)^{f(x)}|x\rangle$$

and the oracle's matrix will be a diagonal matrix of the form:

$$O = \begin{bmatrix} (-1)^{f(0)} & 0 & \dots & 0 \\ 0 & (-1)^{f(1)} & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & (-1)^{f(N-1)} \end{bmatrix}$$

□ The Grover Iteration - Geometric Visualization

* Exactly 1 solution

Grover's algorithm depends on a key concept known as the Grover iteration, G . Each application of G will ask the oracle how close the current state is to a solution.

$$\langle p | \sum_{j \neq p} \frac{1}{\sqrt{N-1}} = \langle p |$$

Each object is encoded in integers $1, \dots, N$, which we'll write equivalently as $0, 1, \dots, N-1$.

We do this to encode the integers in binary $x \in \{0, 1\}^n$, where $n = \log_2 N$ is integral by assumption that N is a power of 2.

So, our solution will be some bit string z , which we don't know.

We have no idea what the solution z is other than it is one of the bit strings $x \in \{0, 1\}^n$. As far as we are concerned, every bit string is equally probable to be the solution. We should thus start in the equal superposition state,

$$|\Phi\rangle := \frac{1}{\sqrt{2^n}} \sum_{x \in \{0, 1\}^n} |x\rangle$$

where,

$$H^{\otimes n} |0\rangle^{\otimes n} = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle = |\Phi\rangle$$

Since $|\Phi\rangle$ is an equal superposition over all bit strings, it includes the correct bit string, say z .

$$|\Phi\rangle = \frac{1}{\sqrt{2^n}} \left(|wrong\rangle + |wrong\rangle + \dots + |wrong\rangle + |z\rangle + |wrong\rangle + \dots + |wrong\rangle \right)$$

If we measure in the computational basis now, we would have merely $\frac{1}{2^n}$ probability of getting the correct bit string, which is very small for large n . In other words, we'd just be guessing!

We want some way to increase, or amplify, this probability. This is the goal of the Grover iteration, which achieves this in a clever way.

The overlap of the correct string $|z\rangle$ with our initial guess is exponentially small in the # of qubits n :

$$\langle z | \phi \rangle = \frac{1}{\sqrt{2^n}}$$

$\Rightarrow$ Geometrically, the state $|\phi\rangle$ is nearly orthogonal to $|z\rangle$. It is exactly orthogonal in the limiting case $n \rightarrow \infty$, and for all practical purposes, for large enough n we can consider them to be essentially orthogonal.

- If the two states $|\phi\rangle$ and $|z\rangle$ were exactly orthogonal, Grover's algorithm would fail.

The angle b/w $|\phi\rangle$ and the vector exactly orthogonal to $|z\rangle$ is,

$$\Delta := \frac{\pi}{2} - \cos^{-1} \frac{1}{\sqrt{2^n}} = \sin^{-1} \frac{1}{\sqrt{2^n}} \approx \frac{1}{\sqrt{2^n}}$$

where we used the small angle approximation which is valid for large n .

The Grover iteration, geometrically, consists of the following steps:

- ① Reflect about the correct solution $|z\rangle$.

$$R_z := 2|z\rangle\langle z| - I$$

- ② Reflect about the equal superposition state (initial guess) $|\phi\rangle$.

$$R_\phi := 2|\phi\rangle\langle\phi| - I$$

These 2 reflections rotate any state vector $|\psi\rangle$ closer to the correct state $|z\rangle$.

Consider the effect of the Grover iteration $G := R_\phi R_z$ on the initial state $|\phi\rangle$.

$$\begin{aligned} R_z|\phi\rangle &= (2|z\rangle\langle z| - I)|\phi\rangle \\ &= \frac{2}{\sqrt{2^n}}|z\rangle - |\phi\rangle \end{aligned} \quad \left(\langle z|\phi\rangle = \frac{1}{\sqrt{2^n}} \right)$$

$$\begin{aligned} R_\phi R_z|\phi\rangle &= R_\phi \left(\frac{2}{\sqrt{2^n}}|z\rangle - |\phi\rangle \right) \\ &= (2|\phi\rangle\langle\phi| - I) \left(\frac{2}{\sqrt{2^n}}|z\rangle - |\phi\rangle \right) \\ &= \left(\frac{4}{2^n} - 1 \right) |\phi\rangle - \frac{2}{\sqrt{2^n}}|z\rangle \end{aligned}$$

$$\sin \Delta = \frac{1}{\sqrt{2^n}}$$

∴
The Grover iteration $G_1 := R_\phi R_z$ has the effect,

$$G_1 |\phi\rangle = (4\sin^2 \Delta - 1) |\phi\rangle - 2\sin \Delta |z\rangle$$

where $N = 2^n$ and n is the # of qubits.

⇒ the amplitude of $|z\rangle$ is increased after applying the Grover iteration.

* The probability of measuring $|z\rangle$ before the Grover iteration,

$$P = |\langle z | \Phi \rangle|^2 = \frac{1}{2^n}$$

The probability of measuring $|z\rangle$ after the Grover iteration,

$$\begin{aligned} P_G &= |\langle z | G | \Phi \rangle|^2 \\ &= |(4 \sin^2 \Delta - 1) \langle z | \Phi \rangle - 2 \sin \Delta \langle z | z \rangle|^2 \\ &= \left| (4 \sin^2 \Delta - 1) \frac{1}{\sqrt{2^n}} - 2 \sin \Delta \right|^2 \\ &= \frac{1}{2^n} (16 \sin^4 \Delta - 8 \sin^2 \Delta + 1) + 4 \sin^2 \Delta - \frac{4 \sin \Delta}{\sqrt{2^n}} (4 \sin^2 \Delta - 1) \\ &= \frac{1}{2^n} + \frac{16 \sin^4 \Delta}{2^n} - \frac{16}{\sqrt{2^n}} \sin^3 \Delta + \left(4 - \frac{8}{2^n}\right) \sin^2 \Delta + \frac{4 \sin \Delta}{\sqrt{2^n}} \\ &= \frac{1}{2^n} + \frac{16 \sin^4 \Delta}{2^n} + \left(4 - \frac{8}{2^n}\right) \sin^2 \Delta - \frac{4}{\sqrt{2^n}} \sin \Delta (4 \sin^2 \Delta - 1) \\ &= \frac{1}{2^n} + \frac{16 \sin^4 \Delta}{2^n} + \left(4 - \frac{8}{2^n}\right) \sin^2 \Delta + \frac{4 \sin \Delta}{\sqrt{2^n}} (1 - 4 \sin^2 \Delta) \end{aligned}$$

$$\sin \Delta = \frac{1}{\sqrt{2^n}} \quad \text{and} \quad n \rightarrow \infty$$

$$4 - \frac{8}{2^n} > 0 \Rightarrow \left(4 - \frac{8}{2^n}\right) \sin^2 \Delta > 0$$

$$\frac{4 \sin \Delta}{\sqrt{2^n}} \left(1 - 4 \sin^2 \Delta\right) = \frac{4}{2^n} \left(1 - \frac{4}{2^n}\right) > 0$$

$$\Rightarrow P_G = \frac{|\langle z|G|\Phi \rangle|^2}{2^n} = \frac{|\langle z|\Phi \rangle|^2}{2^n}$$

* $G|\phi\rangle$ and $-G|\phi\rangle$ produce the same probability distribution and are thus the same state upto global phase.

* The angle θ b/w $|\phi\rangle$ and $-G|\phi\rangle$ is such that

$$\cos\theta = \langle\phi|(-G|\phi\rangle) = 1 - 2\sin^2\Delta = \cos 2\Delta$$

Proof.

$$-G|\phi\rangle = -(4\sin^2\Delta - 1)|\phi\rangle + 2\sin\Delta|z\rangle$$

$$\langle\phi|-G|\phi\rangle = -(4\sin^2\Delta - 1) + 2\sin\Delta \langle\phi|z\rangle$$

$$= -4\sin^2\Delta + 1 + \frac{2\sin\Delta}{\sqrt{2^n}}$$

$$\sin\Delta = \frac{1}{\sqrt{2^n}}$$

$$= -4\sin^2\Delta + 1 + 2\sin^2\Delta = \underline{\underline{1 - 2\sin^2\Delta}}$$

$\therefore$
The angle b/w $|\phi\rangle$ and $-G|\phi\rangle$ is 2Δ .

Here is the eureka! moment.

We started with a vector $|\phi\rangle$ that had angle Δ relative to $|z^\perp\rangle$. After applying the Grover iteration we ended up with a vector $-G|\phi\rangle$ which has angle 2Δ relative to $|\phi\rangle$.

$\Rightarrow -G|\phi\rangle$ has angle 3Δ relative to $|z^\perp\rangle$

$\therefore$
The effect of the Grover iteration is to move the initial vector an angle 2Δ closer to the correct vector.

The Grover iteration G has this effect on any vector, not just $|\phi\rangle$, as long as that vector is not $|z^\perp\rangle$ upto global phase.

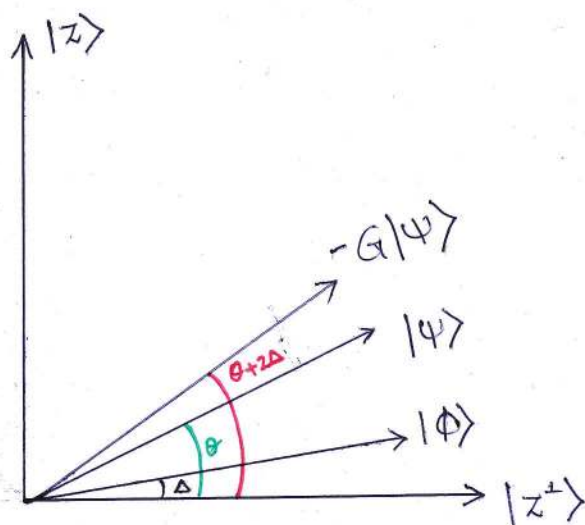
- Grover's algorithm is simply $G^k |\phi\rangle$ for some integer k .

Effect of Grover iteration

Let $|\psi\rangle$ be any state such that $\langle z^\perp | \psi \rangle =: \cos \theta \neq 0$.
Then, the angle b/w $|z\rangle$ and $G|\psi\rangle$ (modulo a global phase) is :

$$\cos^{-1} \langle z | G|\psi \rangle = \theta + 2\Delta$$

$$\text{where } \Delta := \sin^{-1} \frac{1}{\sqrt{N}} \approx \frac{1}{\sqrt{N}}$$



The Quantum Complexity

Grover's algorithm is simply $G^k |\Phi\rangle$ for some integer k . The # of times we need to implement G is precisely the (oracle) complexity of the algorithm.

We want our state to be as close to the vector $|z\rangle$ as possible.

The angle b/w $G^k |\Phi\rangle$ and $|z\rangle$ (i.e., our final state after k iterations of G) is :

$$\theta_k := \Delta + 2k\Delta = (2k+1)\Delta$$

We would like this angle θ_k to be $\frac{\pi}{2} - \Delta$.

Setting $\theta_k = \frac{\pi}{2} - \Delta$,

$$(2k+1)\Delta = \frac{\pi}{2} - \Delta$$

$$\Rightarrow 2(k+1)\Delta = \frac{\pi}{2}$$

$$\therefore k = \frac{\frac{\pi}{2}}{2\Delta} - 1$$

$$\Delta := \sin^{-1}\left(\frac{1}{\sqrt{N}}\right) \approx \frac{1}{\sqrt{N}}$$

$$k \approx \frac{\frac{\pi}{2}}{2} \sqrt{N} - 1$$

We want k to be an integer because we're going to apply G an integer # of times.

So we round k to an integer. However, the farthest k can be from an integer is one-half.

Let θ_k^* be the optimal angle, and let θ_k be the angle we rotate by after rounding k^* to the closest integer k .

$$\downarrow \quad |k - k^*| \leq \frac{1}{2}$$

$$|\theta_k - \theta_k^*| = 2|k - k^*| \Delta \leq \Delta$$

$\therefore$ The effect of rounding error is:

$$\theta_k = \theta_k^* \pm \Delta$$

$\Rightarrow$ We are at worst an angle Δ away from $|z\rangle$ at the end of Grover's algorithm.

Angle b/w $|z\rangle$ and $G^k|\phi\rangle$ is Δ .

For the worst case when the final angle is Δ away from optimal, the probability of success is :

$$P(\text{measure } |z\rangle) = |\langle z | G^k | \phi \rangle|^2 \\ = \cos^2 \Delta = 1 - \sin^2 \Delta = 1 - \frac{1}{N}$$

The Complexity of Grover's algorithm

For large N , after

$$k := \text{round} \left(\frac{\pi}{4} \sqrt{N} - 1 \right)$$

applications of the Grover's iteration $G_i = R_\phi R_z$ to the initial state $|\Phi\rangle$, the probability of measuring the correct solution is

$$P(\text{measure } |x\rangle) = 1 - \frac{1}{N}$$

* M solutions

The equal superposition state is,

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle$$

Let's define normalized states,

$$|\alpha\rangle = \frac{1}{\sqrt{N-M}} \sum_x'' |x\rangle$$

$$|\beta\rangle = \frac{1}{\sqrt{M}} \sum_x' |x\rangle$$

where $\sum_x''$ indicates a sum over all x which are not solutions to the search problem, and $\sum_x'$ indicates a sum over all x which are solutions to the search problem.

The initial state $|\psi\rangle$ may be re-expressed as:

$$|\psi\rangle = \sqrt{\frac{N-M}{N}} |\alpha\rangle + \sqrt{\frac{M}{N}} |\beta\rangle$$

So the initial state of the quantum computer is in the space spanned by $|\alpha\rangle$ and $|\beta\rangle$.

The Grover iteration can be defined by the operator,

$$G = (2|\psi\rangle\langle\psi| - I)O$$

where the oracle operation O performs a reflection about the vector $|\alpha\rangle$ in the plane defined by $|\alpha\rangle$ and $|\beta\rangle$. That is,

$$O(a|\alpha\rangle + b|\beta\rangle) = a|\alpha\rangle - b|\beta\rangle$$

Similarly,

$2|\psi\rangle\langle\psi| - I$ also performs a reflection in the plane defined by $|\alpha\rangle$ and $|\beta\rangle$, about the vector $|\psi\rangle$.

$$\begin{aligned} \text{Ref}(\theta) \text{Ref}(\phi) &= \begin{bmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{bmatrix} \begin{bmatrix} \cos 2\phi & \sin 2\phi \\ \sin 2\phi & -\cos 2\phi \end{bmatrix} \\ &= \begin{bmatrix} \cos[2(\theta-\phi)] & -\sin[2(\theta-\phi)] \\ \sin[2(\theta-\phi)] & \cos[2(\theta-\phi)] \end{bmatrix} = \text{Rot}[2(\theta-\phi)] \end{aligned}$$

$\Rightarrow$ The product of two reflections is a rotation.

$\therefore$ The state $G^k|\psi\rangle$ remains in the space spanned by $|\alpha\rangle$ and $|\beta\rangle$ for all k .

$$\text{Let } \cos \theta/2 = \langle \alpha | \psi \rangle = \sqrt{\frac{N-M}{N}}$$

$$\Rightarrow |\psi\rangle = \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$O|\psi\rangle = \cos \theta/2 |\alpha\rangle - \sin \theta/2 |\beta\rangle$$

$$G|\psi\rangle = (2|\psi\rangle\langle\psi| - I)O|\psi\rangle$$

$$\left(\begin{array}{l} \langle \psi | \alpha \rangle = \cos \theta/2 \quad \& \quad \langle \psi | \beta \rangle = \sin \theta/2 \end{array} \right.$$

$$\rightarrow = (2|\psi\rangle\langle\psi| - I)(\cos \theta/2 |\alpha\rangle - \sin \theta/2 |\beta\rangle)$$

$$= 2\cos^2 \theta/2 |\psi\rangle - 2\sin^2 \theta/2 |\psi\rangle - \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$= 2\cos \theta |\psi\rangle - \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$= 2\cos \theta \cos \theta/2 |\alpha\rangle + 2\cos \theta \sin \theta/2 |\beta\rangle - \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$= \cos \theta/2 (2 \cos \theta - 1) |\alpha\rangle + \sin \theta/2 (2 \cos \theta + 1) |\beta\rangle$$

$$= \cos \theta/2 (4 \cos^2 \theta/2 - 3) |\alpha\rangle + \sin \theta/2 (-4 \sin^2 \theta/2 + 3) |\beta\rangle$$

$$= (4 \cos^3 \theta/2 - 3 \cos \theta/2) |\alpha\rangle + (3 \sin \theta/2 - 4 \sin^3 \theta/2) |\beta\rangle$$

$$\sin 3\alpha = 3 \sin \alpha - 4 \sin^3 \alpha$$

$$\cos 3\alpha = 4 \cos^3 \alpha - 3 \cos \alpha$$

$$\rightarrow G|\psi\rangle = \cos \frac{3\theta}{2} |\alpha\rangle + \sin \frac{3\theta}{2} |\beta\rangle$$

$\Rightarrow$ The rotation angle is in fact θ .

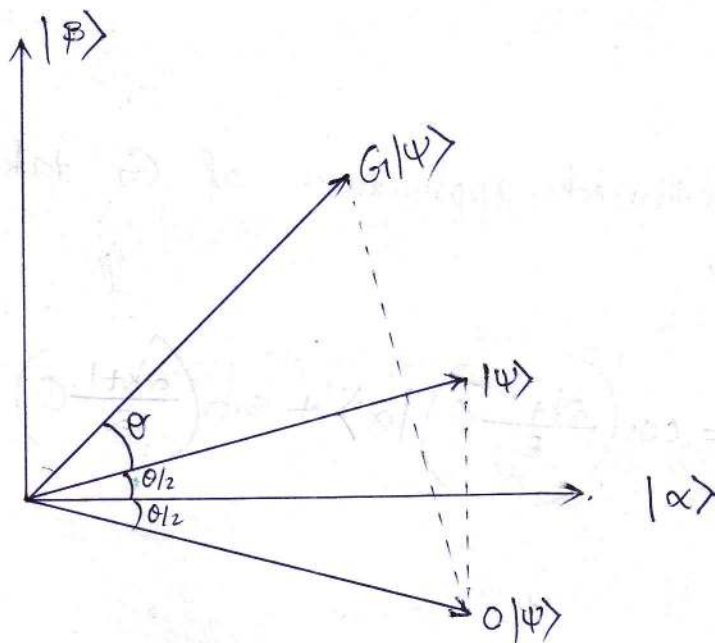
The continued application of G takes the state to,

$$G^k |\psi\rangle = \cos \left(\frac{(2k+1)\theta}{2} \right) |\alpha\rangle + \sin \left(\frac{(2k+1)\theta}{2} \right) |\beta\rangle$$

G is a rotation in the 2D space spanned by $|\alpha\rangle$ and $|\beta\rangle$, rotating the space by θ radians per application of G .

Repeated application of the Grover iteration rotates the state vector close to $|\beta\rangle$.

When this occurs, an observation in the computational basis produces with high probability one of the outcomes superposed in $|\beta\rangle$, that is, a solution to the search problem!



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* In the $|\alpha\rangle, |\beta\rangle$ basis the Grover iteration can be written as,

$$G = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix}$$

where θ is a real # in the range 0 to $\pi/2$ (assuming for simplicity that $M \leq N/2$); this limitation will be lifted shortly), chosen so that

$$\sin\theta = \frac{2\sqrt{M(N-M)}}{N}$$

$$\rightarrow \cos\theta/2 = \langle\alpha|\psi\rangle = \sqrt{\frac{N-M}{N}}$$

$$\Rightarrow \sin\theta/2 = \sqrt{1 - \frac{N-M}{N}} = \sqrt{\frac{M}{N}}$$

$$\sin\theta = 2\sin\theta/2 \cos\theta/2 = 2 \cdot \sqrt{\frac{M}{N}} \sqrt{\frac{N-M}{N}} = \frac{2\sqrt{M(N-M)}}{N}$$

Assuming that $M \leq N/2 \Rightarrow \frac{M}{N} \leq \frac{1}{2} \Rightarrow \sin\theta/2 \leq \frac{1}{\sqrt{2}}$.

$$\Rightarrow \theta/2 \leq \pi/4 \Rightarrow \underline{\underline{\theta \leq \frac{\pi}{2}}}$$

Performance

How many times must the Grover iteration be repeated in order to rotate $|\psi\rangle$ near $|\beta\rangle$?

The initial state of the system is,

$$|\psi\rangle = \sqrt{\frac{N-M}{N}} |\alpha\rangle + \sqrt{\frac{M}{N}} |\beta\rangle$$

$$= \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$\frac{\pi}{2} - \frac{\theta}{2} = \frac{\pi}{2} - \sin^{-1} \sqrt{M/N} = \cos^{-1} \sqrt{M/N}$$

$\therefore$ Rotating thro' $\arccos \sqrt{M/N}$ radians takes the system to $|\beta\rangle$.

Let $CI(x)$ denote the integer closest to the real number x , where by convention we round halves down; Example: $CI(3.5) = 3$.

$$\sin \theta/2 = \sqrt{M/N} \leq 1/\sqrt{2} \implies \theta/2 \leq \pi/4.$$

$$|x - CI(x)| \leq 1/2 \implies \text{angle}, \theta \leq \theta/2 \leq \pi/4$$

$\therefore$ Repeating the Grover iteration

$$R = CI\left(\frac{\arccos \sqrt{M/N}}{\theta}\right)$$

times rotates $|\psi\rangle$ to within an angle $\theta/2 \leq \pi/4$ of $|\Phi\rangle$.

$$\theta/2 \leq \pi/4 \Rightarrow \cos \theta/2 \geq 1/\sqrt{2}$$

$$P(\text{measure } |B\rangle) \geq \cos^2(\theta/2) = 1 - \sin^2 \theta/2 \geq \frac{1}{2} \\ = 1 - \frac{M}{N}$$

Observation of the state in the computational basis then yields a solution to the search problem with probability at least $1/2$.

For specific values of M and N it is possible to achieve a much higher probability of success.

For example,

when $M \leq N$

$$\theta \approx \sin \theta = \frac{2\sqrt{M(N-M)}}{N} \approx 2\sqrt{\frac{M}{N}}$$

The angular error in the final state is at most $\frac{\theta}{2} \approx \sqrt{\frac{M}{N}}$,

giving a probability of error of at most M/N .

$$R = \frac{C}{\theta} \left(\frac{\arccos \sqrt{M/N}}{\theta} \right)$$

Note that,

R depends on the # of solutions M but not on the identity of those solutions, so provided we know M we can apply the quantum search algorithm.

We can even remove the need for a knowledge of M in applying the search algorithm, which will be explained in the section 6.3- Quantum Counting.

$$R = CI \left(\frac{\arccos \sqrt{M/N}}{0} \right)$$

— (6.15)

The form (6.15) is useful as an exact expression for the # of oracle calls used to perform the quantum search algorithm, but it would be useful to have a simpler expression summarizing the essential behavior of R .

The ceiling function maps x to the least integer greater than or equal to x , denoted $\lceil x \rceil$ or $\text{ceil}(x)$.

Ex:- $\lceil 2.4 \rceil = 3$, $\lceil -2.4 \rceil = -2$

$$\frac{\pi}{2} - \frac{\theta}{2} = \frac{\pi}{2} - \sin^{-1} \sqrt{M/N} = \cos^{-1} \sqrt{M/N}$$

$$\Rightarrow \cos^{-1} \sqrt{M/N} \leq \pi/2$$

$$\therefore R = CI \left(\frac{\arccos \sqrt{M/N}}{\theta} \right) \leq \left\lceil \frac{\pi/2}{\theta} \right\rceil$$

$$\Rightarrow R \leq \left\lceil \frac{\pi}{2\theta} \right\rceil$$

So a lower bound on θ will give an upper bound on R .

$$f(x) = \sin x - x, \quad 0 \leq x \leq \pi/2$$

$$f'(x) = \cos x - 1 \Rightarrow \begin{cases} f'(0) = 1 - 1 = 0 \\ f'(\pi/2) = 0 - 1 = -1 \end{cases}$$

$$f''(x) = -\sin x < 0 \quad \left. \vphantom{f''(x)} \right\} f(x) \text{ is a decreasing function in } [0, \pi/2]$$

$$\left. \begin{aligned} f(0) &= 0 \\ f(\pi/2) &= 1 - \pi/2 < 0 \end{aligned} \right\} f(x) = \sin x - x \leq 0$$

$$\sin x \leq x \Rightarrow \frac{\sin x}{x} \leq 1$$

Lemma

Assume that $M \leq N/2$,

$$\left[\begin{array}{l} \sqrt{\frac{M}{N}} \leq \frac{1}{\sqrt{2}} \Rightarrow \sin \frac{\theta}{2} \leq \frac{1}{\sqrt{2}} \\ \frac{\theta}{2} \leq \pi/4 \Rightarrow \theta \leq \pi/2. \end{array} \right.$$

Lemma $\Rightarrow \frac{\theta}{2} \geq \sin \frac{\theta}{2} = \sqrt{M/N}$

$$\Rightarrow \underline{\theta \geq 2\sqrt{M/N}}$$

from which we obtain an elegant upper bound on the # of iterations required,

$$R \leq \left\lceil \frac{\pi}{4} \sqrt{\frac{N}{M}} \right\rceil$$

$\Rightarrow R = O(\sqrt{N/M})$ Grover iterations (and thus oracle calls) must be performed in order to obtain a solution to the search problem with high probability, a quadratic improvement over the $O(N/M)$ ~~oracle~~

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□ The Procedure

The algorithm makes use of a single n qubit register.

The goal of the algorithm is to find a solution to the search problem, using the smallest possible # of applications of the oracle.

The oracle is a unitary operator O , defined by its action on the computational basis:

$$|x\rangle|q\rangle \xrightarrow{O} |x\rangle|q \oplus f(x)\rangle$$

where, $|q\rangle$: oracle qubit, which is flipped if $f(x)=1$, and is unchanged otherwise.

The function f takes as input an integer x , in the range 0 to $N-1$ such that

$f(x) = 1$ if x is a solution to the search problem.

$f(x) = 0$ if x is not a solution to the search problem.

The action of the oracle is,

$$\begin{aligned} |x\rangle \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) &\xrightarrow{O} |x\rangle \left(\frac{|0\rangle \oplus f(x) - |1\rangle \oplus f(x)}{\sqrt{2}} \right) \\ &= |x\rangle \left(\frac{|f(x)\rangle - |\bar{f}(x)\rangle}{\sqrt{2}} \right) \\ &= (-1)^{f(x)} |x\rangle \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \end{aligned}$$

The state of the oracle qubit remains $\frac{|0\rangle - |1\rangle}{\sqrt{2}}$ throughout the quantum search algorithm, & therefore be omitted from further discussions.

with this convention, the action of the oracle may be written as:

$$|x\rangle \xrightarrow{O} (-1)^{f(x)} |x\rangle$$

⇒ The oracle marks the solutions to the search problem, by shifting the phase of the solution.

$$\langle u | \sum_x \frac{1}{M} = \langle u |$$

$$\langle u | \frac{1}{M} = \langle u | \frac{1}{M} = \langle u |$$

$$\langle u | \sum_x \frac{1}{M} = \langle u |$$

$$\langle u | \sum_x \frac{1}{M} = \langle u |$$

$$\frac{1}{\sqrt{2}}$$

&

ms.

The algorithm begins with the computer in the state $|0\rangle^n$. The Hadamard transform is used to put the computer in the equal superposition state,

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle$$

which can be expressed as,

$$|\psi\rangle = \sqrt{\frac{N-M}{N}} |\alpha\rangle + \sqrt{\frac{M}{N}} |\beta\rangle$$

where, $|\alpha\rangle = \frac{1}{\sqrt{N-M}} \sum_x'' |x\rangle$

$$|\beta\rangle = \frac{1}{\sqrt{M}} \sum_x' |x\rangle$$

and $\sum_x'$ indicates a sum over all x which are solutions to the search problem, and $\sum_x''$ indicates a sum over all x which are not

solutions to the search problem.

So the initial state of the quantum computer is in the space spanned by $|\alpha\rangle$ and $|\beta\rangle$.

The quantum search algorithm consists of repeated application of a quantum subroutine, known as the Grover iteration or Grover's iteration, G .

$$G = (2|\psi\rangle\langle\psi| - I)O$$

$$= H^{\otimes n} (2|0\rangle\langle 0| - I) H^{\otimes n} O$$

where, the oracle operation O performs a reflection about the vector $|\alpha\rangle$ in the plane defined by $|\alpha\rangle$ and $|\beta\rangle$.

$$\text{i.e., } O(a|\alpha\rangle + b|\beta\rangle) = a|\alpha\rangle - b|\beta\rangle$$

* $2|\psi\rangle\langle\psi| - I$ also performs a reflection in the plane defined by $|\alpha\rangle$ and $|\beta\rangle$, about the vector $|\psi\rangle$.

* The Householder matrix for a reflection about the hyperplane $\perp$ to a unit vector $\hat{u}$ is:

$$H = I - 2\hat{u}\hat{u}^\dagger$$

The key insight to implementing G is actually the intuitive interpretation of reflections:

Do nothing to the $\parallel$ component, and add a minus sign to the $\perp$ component.

Fig:- A quantum circuit that implements a reflection about $|0\rangle^{\otimes n}$, i.e., the unitary $R_0 = I - 2|0\rangle\langle 0|$. Here $n=3$, and a single ancilla is in the $|-\rangle$ state is utilized for any n . The multicontrolled Toffoli gate can be constructed from standard Toffoli gates.

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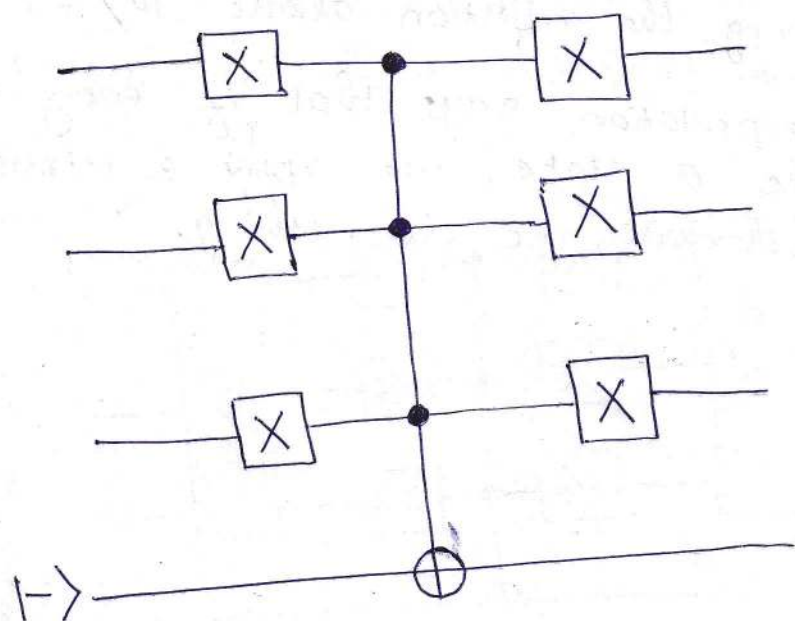
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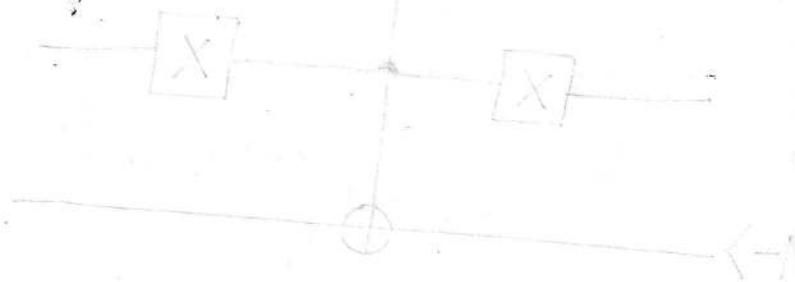


$$U = I - 2|X\rangle\langle X| = I - 2|X\rangle\langle X|$$

where V is a unitary operator from the ground state to the excited state.

Considering the reflection about $|0\rangle^{\otimes n}$.

The interpretation says that if every qubit is in the 0 state, we apply a minus sign. Otherwise, we do nothing.



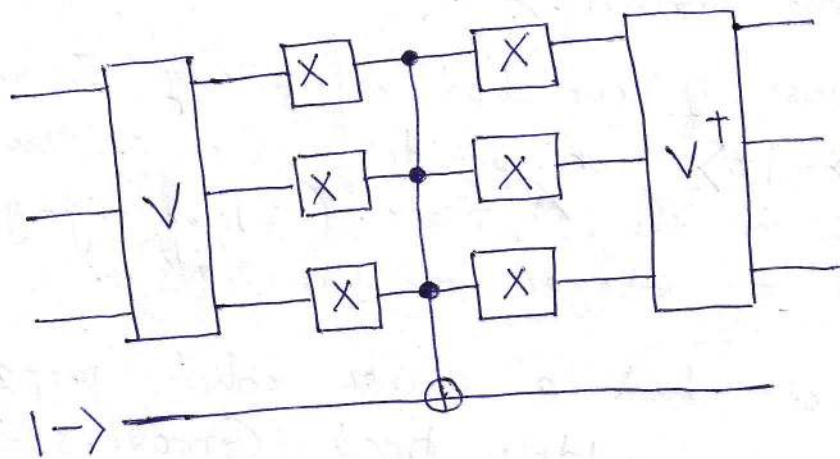
We can modify the reflection about $|0\rangle$, to reflect about any state $|v\rangle$.

$$R_v = I - 2|v\rangle\langle v| = I - 2V|0\rangle\langle 0|V^\dagger$$

where V is a unitary operator which prepares $|v\rangle$ from the ground state.
i.e., $V|0\rangle = |v\rangle$.

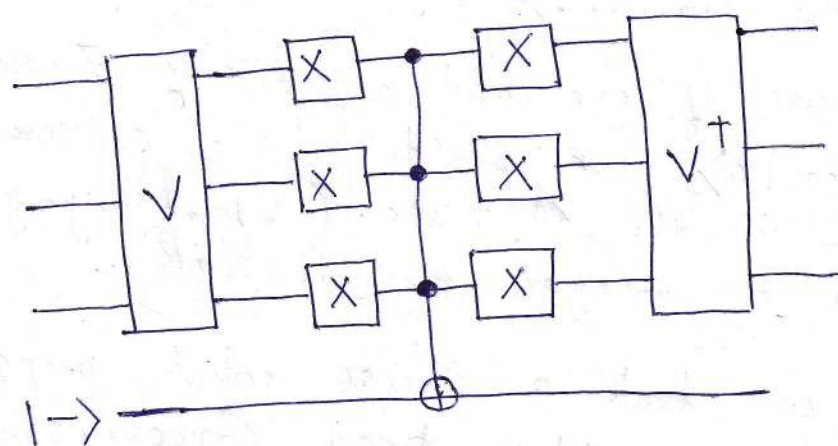
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Fig:- A quantum circuit for performing the reflection $R_V = I - 2|v\rangle\langle v|$ where $|v\rangle = V|0\rangle$



ubit

Fig:- A quantum circuit for performing the reflection $R_v = I - 2|v\rangle\langle v|$ where $|v\rangle = V|0\rangle$



10/11 - We can not implement the reflection about the solution state $|z\rangle$ in the same manner as in the circuit.

Because, if we had a unitary Z such that $Z|0\rangle = |z\rangle$, we wouldn't have a problem to solve in the 1st place. (We're trying to find an integer $z \in \{1, 2, \dots, N\}$).

If we had a circuit which prepared $|z\rangle$, we wouldn't need Grover's algorithm!

All we are allowed is the Quantum oracle which acts on a state of $n = \log N$ qubits and sets an ancilla qubit to be $|1\rangle$ if the register is a solution to the search problem. This is precisely what we want to implement the reflection about $|z\rangle$. We simply act with the oracle O on the n register and ancilla to perform the reflection, $R_z = I - 2|z\rangle\langle z|$.

The Grover iteration G may be broken up into 4 steps:

- ① Apply the Oracle O
- ② Apply the Hadamard transform $H^{\otimes n}$
- ③ Perform a conditional phase shift on the computer, with every computational basis except $|0\rangle$, receiving a phase shift of -1 .

$$|x\rangle \rightarrow -(-1)^{\delta_{x0}} |x\rangle$$

- ④ Apply the Hadamard transform $H^{\otimes n}$

- The unitary operator corresponding to the phase shift in the Grover iteration is $2|0X0\rangle - I$.

$$I - 2|0X0\rangle\langle 0X0| = H(I - 2|0X0\rangle\langle 0X0|)H$$

Each of the operations in the Grover iteration may be efficiently implemented on a quantum computer.

The Grover iteration requires only a single oracle call. The combined effect of steps 2, 3 and 4 is:

$$H^{\otimes n} (2|0\rangle\langle 0| - I) H^{\otimes n} = 2|\psi\rangle\langle\psi| - I$$

where $|\psi\rangle$ is the equally weighted superposition of states. Thus the Grover iteration G may be written, $G = (2|\psi\rangle\langle\psi| - I)O$.

Ex: 6.2 Show that the operation $(2|\psi\rangle\langle\psi| - I)$
applied to a general state $\sum_k \alpha_k |k\rangle$

produces

$$\sum_k [-\alpha_k + 2\langle\alpha\rangle] |k\rangle$$

where $\langle\alpha\rangle = \sum_k \frac{\alpha_k}{N}$ is the mean value

of the α_k . For this reason, $(2|\psi\rangle\langle\psi| - I)$
is sometimes referred to as the inversion
about mean operation.

Algorithm: Quantum Search

Inputs: ① A black box oracle O which performs the transformation, $O|x\rangle|y\rangle = |x\rangle|y \oplus f(x)\rangle$ where $f(x) = 0$ for all $0 \leq x < 2^n$ except x_0 , for which $f(x_0) = 1$

② $n+1$ qubits in the state $|0\rangle$

Outputs: x_0

Runtime: $O(\sqrt{2^n})$ operations. Succeeds with probability $O(1)$.

Procedure:

① $|0\rangle^{\otimes n} |0\rangle$

② $\rightarrow \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]$

③ $\left[2|\psi\rangle\langle\psi| - I \right] O \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]$
 $\approx |x_0\rangle \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]$

④ $\rightarrow x_0$

initial state

apply $H^{\otimes n}$ to the 1st n qubits, and HX to the last qubit.

apply the Grover iteration $R \approx \frac{\pi\sqrt{2^n}}{4}$ times.

measure the 1st n qubits.

What happens when more than half the items are solutions to the search problem, i.e., $M \geq N/2$?

$$\sin \theta/2 = \sqrt{\frac{N-M}{N}} \quad \& \quad \cos \theta/2 = \sqrt{\frac{M}{N}}$$

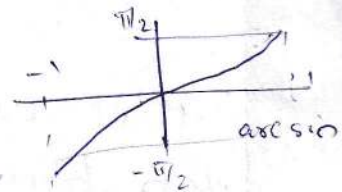
$$\sin \theta = 2 \sin \theta/2 \cos \theta/2$$

$$= 2 \times \sqrt{\frac{N-M}{N}} \times \sqrt{\frac{M}{N}} = \frac{2\sqrt{M(N-M)}}{N}$$

$$\Rightarrow \theta = \arcsin \left(\frac{2\sqrt{M(N-M)}}{N} \right)$$

$$M = N/2 : \theta = \sin^{-1}(1) = \pi/2$$

$$M = N : \theta = \sin^{-1}(0) = 0$$



$\Rightarrow$ The angle θ gets smaller as M varies from $N/2$ to N . As a result, the # of iterations needed by the search algorithm increases with M , for $M \geq N/2$.

Intuitively, this is a silly property for a search algorithm to have: we expect that it should become easier to find a solution to the problem as the # of solutions increases.

② If M is known in advance to be larger than $N/2$.

then we can just randomly pick an item from the search space, and then check that it is a solution using the oracle. This approach has a success probability at least $1/2$, and only requires one consultation with the oracle.

It has the disadvantage that we may not know the # of solutions M in advance.

Intuitively, this is a silly property for a search algorithm to have: we expect that it should become easier to find a solution to the problem as the # of solutions increases.

i) If M is known in advance to be larger than $N/2$.

then we can just randomly pick an item from the search space, and then check that it is a solution using the oracle. This approach has a success probability at least $1/2$, and only requires one consultation with the oracle.

It has the disadvantage that we may not know the # of solutions M in advance.

ii) If it isn't known whether $M \geq N/2$

The idea is to double the # of elements in the search space by adding N extra items to the search space, none of which are solutions. This is effected by adding a single qubit $|2\rangle$ to the search index, doubling the # of items to be searched to $2N$.

A new augmented oracle O' is constructed which marks an item only if it is a solution to the search problem and the extra bit is set to zero. The new search problem has only M solutions out of $2N$ entries. So running the search algorithm with the new oracle O' , at most $R = \frac{\pi}{4} \sqrt{\frac{2N}{M}}$ calls to O' are required, and it follows that $O(\sqrt{\frac{N}{M}})$ calls

to 0 are required to perform the search.

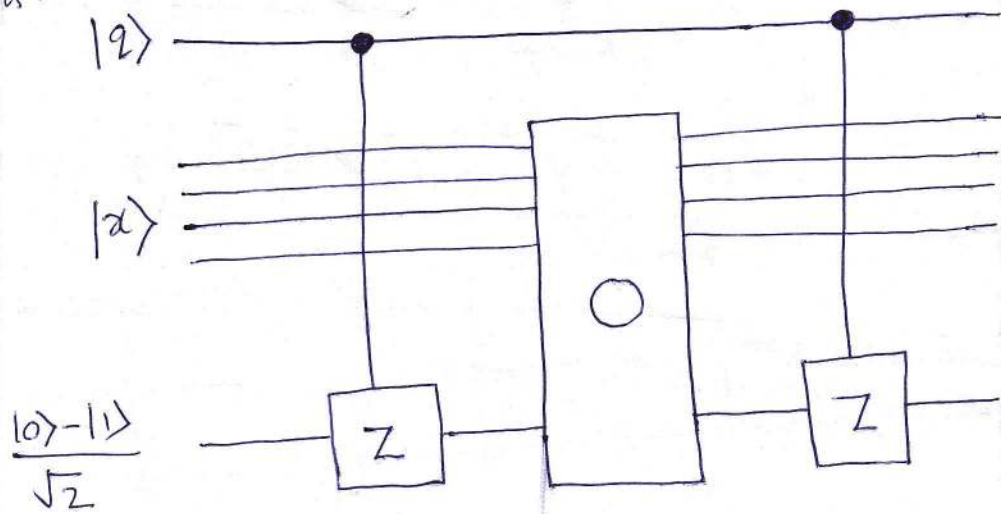
Ex: 6

Ans

search.

Ex: 6.5 Show that the augmented oracle O' may be constructed using one application of O , and elementary quantum gates, using the extra qubit $|2\rangle$.

Ans:



Z gate (Phase flip matrix)

$$Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = |0\rangle\langle 0| - |1\rangle\langle 1|$$

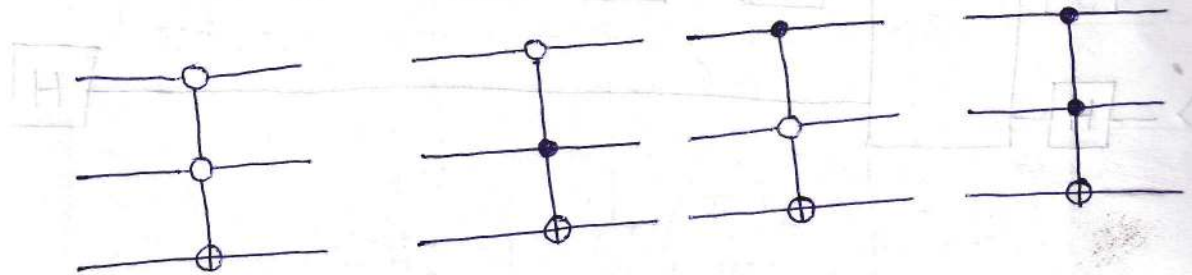
$$\begin{cases} Z|0\rangle = |0\rangle \\ Z|1\rangle = -|1\rangle \end{cases}$$

$$Z \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$$

Quantum Search $\rightarrow$ a two-bit example

Search space size, $N=4$

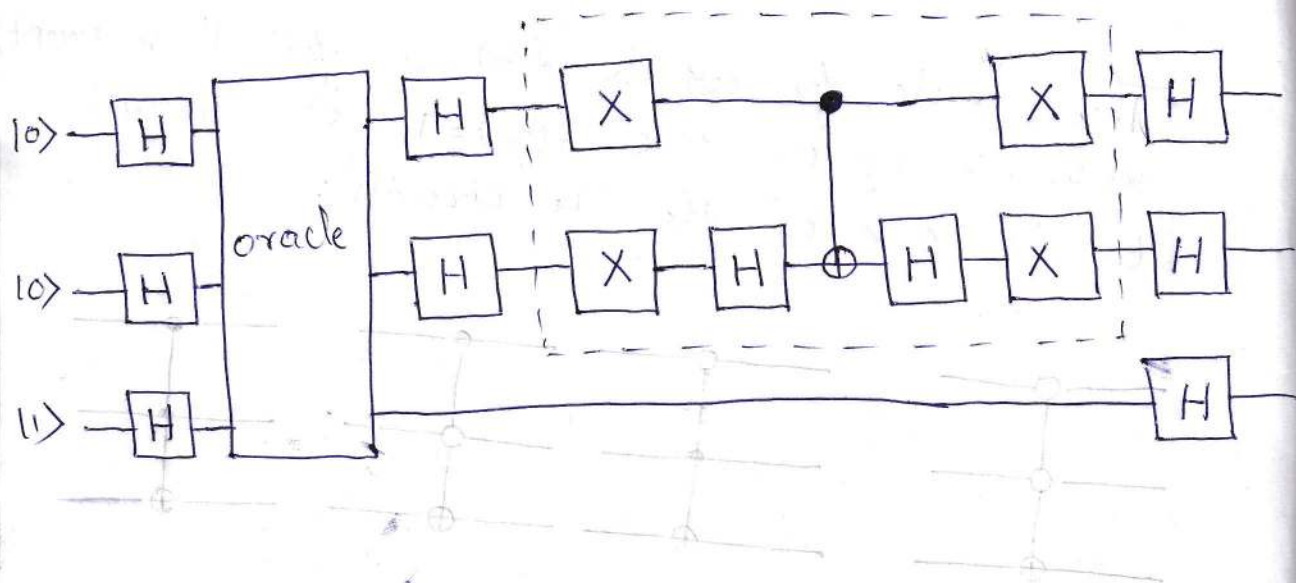
The oracle, for which $f(x)=0$ for all x except $x=x_0$, in which case $f(x_0)=1$, can be taken to be one of the 4 circuits:



corresp. to $x_0=0, 1, 2$, or 3 from left to right, where the top two qubits carry the query x , and the bottom qubit carries the oracle's response.

We need to identify which one among the four.

The quantum circuit which performs the initial Hadamard transforms & a single Grover iteration G is :



The gates in the dotted box performs the conditional phase shift operation $2|00\rangle\langle 00| - I$.

How many times must we repeat G to obtain α_0 ?

$$R = CI \left(\frac{\arccos \sqrt{M/N}}{\theta} \right)$$

Repeating Grover iteration R number of times rotates $|\psi\rangle$ to within an angle $\theta/2 \leq \pi/4$ of $|\beta\rangle$, the superposition of all solutions to the search problem.

$$M=1, N=4$$

$$R = CI \left(\frac{\arccos 1/2}{\theta} \right) = CI \left(\frac{\pi}{3\theta} \right)$$

$$= CI(1) = I$$

$$\left\{ \begin{aligned} \sin \theta &= \frac{2\sqrt{M(N-M)}}{N} \\ &= \frac{2\sqrt{3}}{4} \\ &= \sqrt{3}/2 \\ &\Rightarrow \theta = \pi/3 \end{aligned} \right.$$

$\Rightarrow$ Only exactly one iteration is required to perfectly obtain α_0 .

In the geometric picture, our initial state $|\psi\rangle = \frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle)$ is 30° from $|\alpha\rangle$, and a single rotation by $\theta = 60^\circ$ moves $|\psi\rangle$ to $|\beta\rangle$.

Measurement of the top two qubits gives α_0 after using the oracle only once.

In contrast, a classical computer - or classical circuit - try to differentiate b/w the 4 oracles would require on average 2.25 oracle queries!



Quantum Counting

How quickly can we determine the # of solutions, M , to an N item search problem, if M is not known in advance?

On a classical computer it takes $\mathcal{O}(N)$ consultations with an oracle to determine M .

On a quantum computer it is possible to estimate the # of solutions much more quickly than is possible on a classical computer by combining the Grover iteration with the phase estimation technique based upon the quantum Fourier transform.

This has some important applications:

① If we can estimate the # of solutions quickly then it is also possible to find a solution quickly, even if the # of solutions is unknown, by ^{1st} counting the # of solutions, and then applying the quantum search algorithm to find a solution.

2.29 Quantum counting allows us to decide whether or not a solution even exists, depending on whether the # of solutions is zero, or non-zero. □

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Ex: 6.13 Consider a classical algorithm for the counting problem which samples uniformly and independently, k times from the search space; and let $X_1, \dots, X_k$ be the results of the oracle calls, that is, $X_j = 1$ if the j th oracle call revealed a solution to the problem, and $X_j = 0$ if the j th oracle call did not reveal a solution to the problem. This algorithm returns the estimate $S = N \times \sum_j X_j / k$ for the # of solutions to the search problem. Show that

the standard deviation in S is $\Delta S = \sqrt{M(N-M)/k}$.
 Prove that to obtain a probability at least $3/4$ of estimating M correctly in within an accuracy $\sqrt{M}$ for all values of M we must have $k = \Omega(N)$.

Ans: S : # of solutions to the search problem.
 estimate of the obtained.

$$\text{Var}(S) = E[S^2] - (E[S])^2$$

$$E(S) = E\left(N \times \sum_j \frac{X_j}{k}\right)$$

$$= \frac{N}{k} E\left(\sum_j X_j\right)$$

$$= \frac{N}{k} \sum_j E(X_j)$$

$$= \frac{N}{k} \sum_j \frac{M}{N}$$

$$= \frac{N}{k} \times \frac{kM}{N} = \underline{\underline{M}}$$

Binomial trial & each
done independently.

$$E(X_j) = 1 \times P(X_j=1) + 0 \times P(X_j=0) \\ = P(X_j=1) = \frac{M}{N}$$

S is an unbiased estimator of M .

where we have used:

- ① For a random variable X with a finite list $x_1, \dots, x_k$ of possible outcomes, each of which (respectively) has probability $p_1, \dots, p_k$ of occurring. The expectation value of X is defined as:

$$E[X] = x_1 p_1 + \dots + x_k p_k = \sum_{j=1}^k x_j p_j$$

which can be interpreted as a weighted average of the x_j value, with weights given by their probabilities p_j .

- ② For any random variables X and Y , and a constant 'a', we have

$$E(X+Y) = E(X) + E(Y)$$

$$E(aX) = aE(X)$$

- ③ The expected value of a Binomial distribution is np .

Proof

$$\text{II } P(X=k) = \binom{n}{k} p^k (1-p)^{n-k}$$

$$E[X] = \sum_{k=0}^n k P(X=k)$$

$$= \sum_{k=0}^n k \binom{n}{k} p^k (1-p)^{n-k}$$

$$= \sum_{k=1}^n k \frac{n!}{(n-k)! k!} p^k (1-p)^{n-k}$$

Factoring out np

$$= np \sum_{k=1}^n \frac{(n-1)!}{(n-k)! (k-1)!} p^{k-1} (1-p)^{n-k}$$

$$= np \sum_{k=1}^n \binom{n-1}{k-1} p^{k-1} (1-p)^{n-k}$$

$$= np \sum_{k=1}^{n-1} \binom{n-1}{k-1} p^{k-1} (1-p)^{n-k}$$

$$= np \sum_{j=0}^{n-1} \binom{n-1}{j} p^j (1-p)^{n-1-j} = np$$

where, $j = k-1 \Rightarrow k: 1 \rightarrow n \Rightarrow j: 0 \rightarrow n-1$

$$(x+y)^m = \sum_{j=0}^m \binom{m}{j} x^j y^{m-j}$$

$$\Rightarrow (p+(1-p))^m = 1 = \sum_{j=0}^m \binom{m}{j} p^j (1-p)^{m-j}$$

2 Let $B_i = 1$ if we have a success on the i th trial, and 0 otherwise.

Then the number X of successes is:

$$X = B_1 + B_2 + \dots + B_n.$$

By the linearity of expectation,

$$E(X) = E(B_1 + B_2 + \dots + B_n) = E(B_1) + E(B_2) + \dots + E(B_n)$$

$$E(B_i) = 1 \times P + 0 \times (1-P) = P$$

$$\Rightarrow E(X) = \underbrace{P + P + \dots + P}_{n \text{ times}} = \underline{\underline{nP}}$$

$$E(s^2) = E \left[\left(N \sum_{j=1}^k \frac{X_j}{k} \right)^2 \right]$$

$$= \frac{N^2}{k^2} E \left(\sum_{j=1}^k X_j \right)^2$$

$$= \frac{N^2}{k^2} E \left(\sum_{i=1}^k X_i \times \sum_{j=1}^k X_j \right)$$

$$= \frac{N^2}{k^2} E \left(\sum_{i=1}^k \sum_{j=1}^k X_i X_j \right)$$

$$= \frac{N^2}{k^2} \sum_{i=1}^k \sum_{j=1}^k E(X_i X_j)$$

case 1 : $i = j$

$$E(X_i X_i) = E(X_i^2) = 1^2 \times P(X_i = 1) + 0 \times P(X_i = 0)$$

$$= P(X_i = 1)$$

$$= \frac{M}{N}$$

Case 2: $i \neq j$

$$\begin{aligned} E(X_i X_j) &= 1 \times 1 \times P(X_i=1, X_j=1) + 1 \times 0 \times P(X_i=1, X_j=0) \\ &\quad + 0 \times 1 \times P(X_i=0, X_j=1) + 0 \times 0 \times P(X_i=0, X_j=0) \\ &= P(X_i=1, X_j=1) \\ &= P(X_i=1) P(X_j=1) \\ &= \frac{M}{N} \times \frac{M}{N} = \frac{M^2}{N^2} \end{aligned}$$

Case 1 happens k times $\Rightarrow$ case 2 must happen $k^2 - k$ times.

$$\begin{aligned} \therefore E(S^2) &= \frac{N^2}{k^2} \sum_{i=1}^k \sum_{j=1}^k E(X_i X_j) \\ &= \frac{N^2}{k^2} \left[k \frac{M}{N} + (k^2 - k) \frac{M^2}{N^2} \right] \\ &= \frac{MN}{k} + M^2 - \frac{M^2}{k} \end{aligned}$$

$$\text{Var}(s) = E(s^2) - (E(s))^2$$

$$= \frac{MN}{k} + \cancel{M^2} - \frac{M^2}{k} - \cancel{M^2}$$

$$= \frac{M(N-M)}{k}$$

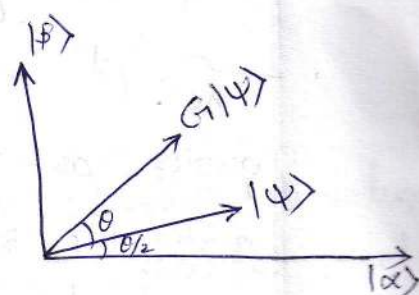
$$\text{s.d.}(s) = \Delta S = \underline{\underline{\sqrt{M(N-M)/k}}}$$

Quantum Counting

Quantum counting is an application of the phase estimation procedure to estimate the eigenvalues of the Grover iteration, G , which in turn enables us to determine the # of solutions M to the search problem.

In the $|\alpha\rangle, |\beta\rangle$ basis:

$$G = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix}$$



Suppose $|a\rangle$ and $|b\rangle$ are the two eigenvectors of the Grover iteration in the space spanned by $|\alpha\rangle$ and $|\beta\rangle$.

Let θ be the angle of rotation determined by the Grover iteration.

→ The corresponding eigenvalues are $e^{i\theta}$ and $e^{i(2\pi-\theta)}$.

For ease of analysis it is convenient to assume that the oracle has been augmented, expanding the size of the search space to $2N$ and ensuring that $\sin^2 \theta/2 = M/2N$.

The function of the phase estimation circuit used for quantum counting is to estimate θ to m bits of accuracy, with a probability of success at least $1-\epsilon$.

The 1st register contains $t = m + \log_2(2 + 1/\epsilon)$ qubits, as per the phase estimation algorithm, and the 2nd register contains $n+1$ qubits, enough to implement the Grover iteration on the augmented search space.

$$N = 2^n \implies 2N = 2^{n+1}$$

The state of the 2nd register is initialized to an equal superposition of all possible inputs $\sum_x |x\rangle$ by a Hadamard transform.

This state is a superposition of the eigenstates $|a\rangle$ and $|b\rangle$.

So the phase estimation circuit gives us an estimate of θ or $2\pi - \theta$ accurate to within $|\Delta\theta| \leq 2^{-m}$, with probability at least $1 - \epsilon$.

Furthermore,

an estimate for $2\pi - \theta$ is clearly equivalent to an estimate of θ with the same level of accuracy, so effectively the phase estimation algorithm determines θ to an accuracy 2^{-m} with probability $1 - \epsilon$.

$\therefore$ Using the equation $\sin^2 \theta/2 = M/2N$ and
 our estimate for θ , we obtain an estimate
 of the # of solutions, M .

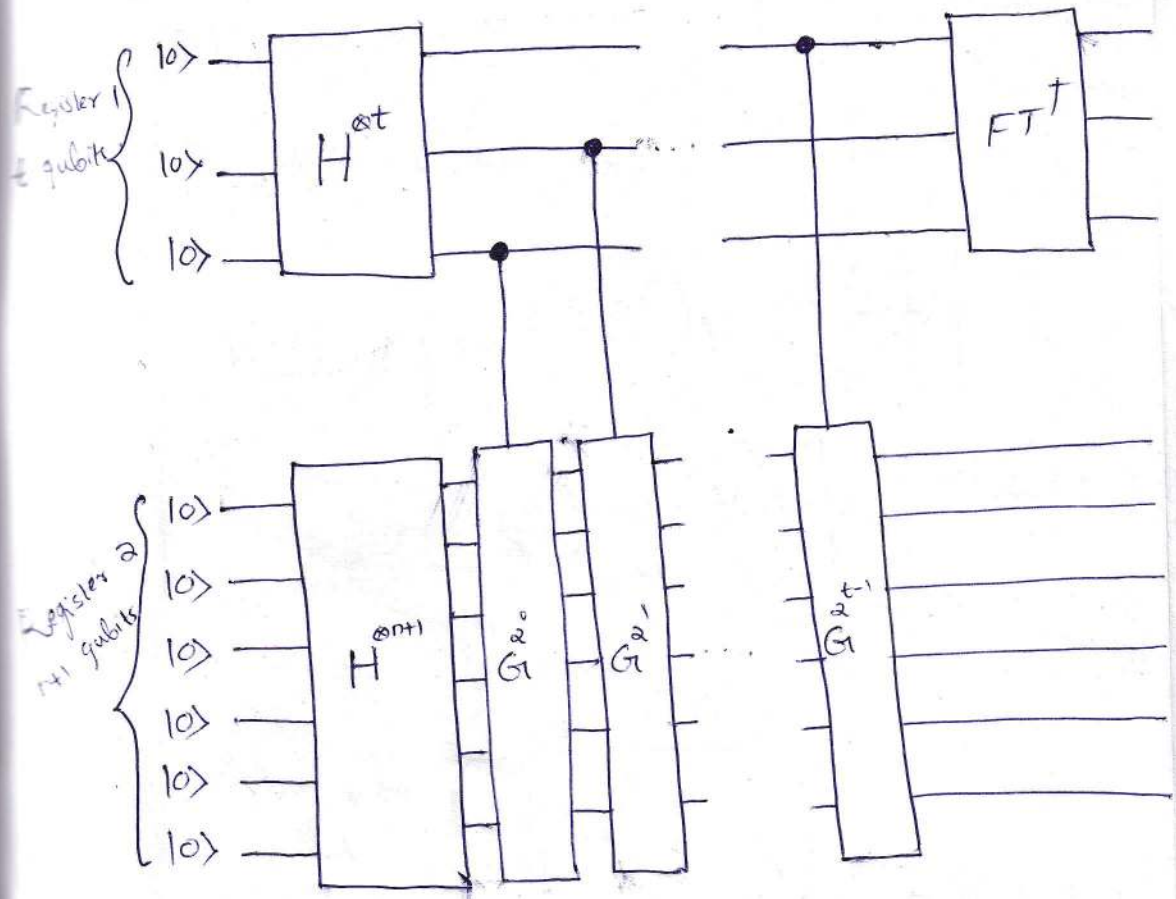
$$\langle \epsilon \rangle \approx \sum$$

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* Circuit for performing approximate quantum counting on a quantum computer.

How large an error, ΔM , is there in the estimate?

$$\sin^2(\theta/2) = M/2N$$

$$\frac{|\Delta M|}{2N} = \left| \sin^2\left(\frac{\theta + \Delta\theta}{2}\right) - \sin^2\left(\frac{\theta}{2}\right) \right|$$

$$= \left(\sin\left(\frac{\theta + \Delta\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right) \right) \left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \right|$$

$$\frac{dy}{dx} = \lim_{\Delta x \rightarrow 0} \frac{f(x + \Delta x) - f(x)}{\Delta x}$$

$$\frac{d}{d\theta} \sin \theta/2 = \cos \theta/2 = \lim_{\Delta\theta/2 \rightarrow 0} \frac{\sin(\frac{\theta + \Delta\theta}{2}) - \sin(\theta/2)}{\Delta\theta/2} \leq 1$$

$$\text{Calculus} \Rightarrow \left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \right| \leq |\Delta\theta|/2$$

$$|a| - |b| \leq |a - b|$$

$$\left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) \right| - \left| \sin\left(\frac{\theta}{2}\right) \right| \leq \left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \right|$$

$$\sin\frac{\theta}{2} = \sqrt{\frac{M}{2N}} > 0$$

$$\left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) \right| - \sin\frac{\theta}{2} \leq \left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \right| \leq \frac{|\Delta\theta|}{2}$$

$$\Rightarrow \left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) \right| < \sin\left(\frac{\theta}{2}\right) + \frac{|\Delta\theta|}{2}$$

$$\therefore \sin\left(\frac{\theta + \Delta\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right) < 2\sin\frac{\theta}{2} + \frac{|\Delta\theta|}{2}$$

and

$$\left| \sin\left(\frac{\theta + \Delta\theta}{2}\right) - \sin\left(\frac{\theta}{2}\right) \right| \leq \frac{|\Delta\theta|}{2}$$

$$\frac{|\Delta M|}{2N} < \left(2 \sin(\theta/2) + \frac{|\Delta \theta|}{2} \right) \frac{|\Delta \theta|}{2}$$

Substituting $\sin^2 \theta/2 = M/2N$ and $|\Delta \theta| \leq 2^{-m}$,

$$\frac{|\Delta M|}{2N} < \left(2 \sqrt{M/2N} + \frac{1}{2^{m+1}} \right) 2^{-m-1}$$

$$|\Delta M| < \left(\sqrt{2MN} + \frac{N}{2^{m+1}} \right) 2^{-m}$$

Example

Suppose we choose $m = \lceil n/2 \rceil + 1$ and $\epsilon = 1/6$.

$$\begin{aligned} t &= m + \log_2 \left(2 + \frac{1}{2\epsilon} \right) \\ &= \lceil n/2 \rceil + 1 + \log_2 (2 + 3) \\ &= \lceil n/2 \rceil + 1 + \log_2 (5) \\ &\approx \lceil n/2 \rceil + 3 \end{aligned} \quad \left. \begin{aligned} \log_2 4 &< \log_2 5 < \log_2 8 \\ \frac{n}{2} &\leq \lceil n/2 \rceil < \frac{n}{2} + 1 \\ \frac{n}{2} + 3 &< t < \frac{n}{2} + 5 \end{aligned} \right\}$$

$$\begin{aligned} \# \{ \text{Grover iterations} \} &= 1 + 2 + 2^2 + \dots + 2^{t-1} \\ &= 2^t - 1 \\ &= 2^{\lceil n/2 \rceil + 1 + \log_2 5} - 1 \end{aligned}$$

$$2^{\frac{n}{2} + 1 + \log_2 4} - 1 < 2^{\lceil n/2 \rceil + 1 + \log_2 5} - 1 < 2^{\frac{n}{2} + 1 + 1 + \log_2 8} - 1$$

$$2^{\frac{n}{2} + 3} - 1 < 2^{\lceil n/2 \rceil + 1 + \log_2 5} - 1 < 2^{\frac{n}{2} + 5} - 1$$

$$2^3 \cdot 2^{n/2} - 1 < 2^{\lceil n/2 \rceil + 1 + \log_2 5} - 1 < 2^5 \cdot 2^{n/2} - 1$$

$$2^3 \sqrt{N} - 1 < 2^{\lceil n/2 \rceil + 1 + \log_2 5} - 1 < 2^5 \sqrt{N} - 1$$

$\Rightarrow$ The algorithm requires $\Theta(\sqrt{N})$ Grover iterations, and thus $\Theta(\sqrt{N})$ oracle calls.

$$|\Delta M| < \left(\sqrt{2MN} + \frac{N}{2^{m+1}} \right) 2^{-m}$$

$$N = 2^n$$

$$\& m = \left\lceil \frac{n}{2} \right\rceil + 1 > \frac{n}{2} + 1 \Rightarrow \frac{n}{2} + 1 - m < 0$$

$$\Rightarrow 2\left(\frac{n}{2} + 1 - m\right) - 1 < -1 < 0$$

$$|\Delta M| < \left(\sqrt{2M \times 2^n} + \frac{2^n}{2^{m+1}} \right) 2^{-m} = \left(\sqrt{\frac{M}{2} \times 2^{n+2}} + 2^{n-m-1} \right) 2^{-m}$$

$$= \left(\sqrt{\frac{M}{2}} \times 2^{\frac{n}{2}+1} + 2^{n-m-1} \right) 2^{-m}$$

$$= \sqrt{\frac{M}{2}} \times 2^{\frac{n}{2}+1-m} + \frac{1}{4} \times 2^{n-2m+1}$$

$$= \sqrt{\frac{M}{2}} \times 2^{\frac{n}{2}+1-m} + \frac{1}{4} \times 2^{2\left(\frac{n}{2}+1-m\right)-1}$$

$$|\Delta M| < \sqrt{\frac{M}{2}} + \frac{1}{4} = \mathcal{O}(\sqrt{M})$$

cover
ells.

The example just described serves double duty as an algorithm for determining whether a solution to the search problem exists at all, i.e., whether $M=0$ or $M \neq 0$.

If $M=0$ we have $|AM| < \frac{1}{4}$, so the algorithm must produce the estimate zero with probability at least $1 - \epsilon = 1 - \frac{1}{6} = \frac{5}{6}$.

Conversely,
if $M \neq 0$ then it is easy to verify that the estimate for M is not equal to 0 with probability at least $\frac{5}{6}$.

Another application of quantum counting is to find a solution to a search problem when the number M of solutions is unknown.

The difficulty in applying the quantum search algorithm is that the # of times to repeat the Grover iteration, depends on knowing the # of solutions M , since

$$R \leq \left\lceil \frac{\pi}{4} \sqrt{N/M} \right\rceil$$

This problem can be alleviated by using the quantum counting algorithm to 1st estimate θ and M to high accuracy using phase estimation, and then to apply the quantum search algorithm, repeating the Grover iteration a # of times determined by $R = \frac{\pi}{\theta} \left(\arccos \sqrt{M/N} \right)$,

with the estimates for θ and M obtained by phase estimation, substituted to determine R .

$$\frac{1001}{9} \cdot \frac{1}{4} = \frac{1001}{36} \cdot \frac{\pi}{4} = \text{round} \left(\frac{1001}{36} \right) \cdot M$$

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with the estimates for θ and M obtained by phase estimation, substituted to determine R .

$$\frac{|\langle \psi | \psi \rangle|}{\theta} = \frac{|\langle \psi | \psi \rangle|}{\theta} \times \frac{\pi}{4} = \frac{\pi}{4} \sqrt{N/M}$$

$$|\psi\rangle = \cos \theta/2 |\alpha\rangle + \sin \theta/2 |\beta\rangle$$

$$\text{where } \cos \theta/2 = \sqrt{\frac{N-M}{N}} \quad \& \quad \sin \theta/2 = \sqrt{\frac{M}{N}}$$

$$\text{Assuming } M \leq N/2 \Rightarrow \frac{M}{N} = \frac{1}{2} \Rightarrow \sqrt{\frac{M}{N}} = \frac{1}{\sqrt{2}} = \sin \theta/2$$

$$0 \leq \theta/2 \leq \pi/4 \Rightarrow 0 \leq \theta \leq \pi/2$$

$$|x - CI(x)| \leq \frac{1}{2} \Rightarrow \text{angle } \theta \leq \frac{\theta}{2} \leq \frac{\pi}{4}$$

$\therefore$ The Grover iterations rotates the initial state $|\psi\rangle$ by θ a few times, such that $|\psi\rangle$ will be rotated to within an angle $\theta/2 \leq \pi/4$ of $|\beta\rangle$.

Using the Quantum Counting algorithm to get an estimate of $\theta \approx \theta + \Delta\theta$ and thus $M \approx M + \Delta M$ which is the # of solutions.

The corresponding error to our initial state $|\psi\rangle$ is $\Delta\theta/2$.

Since, $\frac{\theta}{2} \leq \frac{\pi}{4}$, the maximum error tied to the initial state $|\psi\rangle$ is,

$$\text{Max. Initial Error} = \frac{\pi}{4} \times \frac{|\Delta\theta/2|}{\theta/2} = \frac{\pi}{4} \times \frac{|\Delta\theta|}{\theta}$$

Grover iterations will rotate $|\psi\rangle$ close to $|\beta\rangle$ with the maximum angular separation being $\pi/4$ irrespective of the value of the initial angle which is $\frac{\theta + \Delta\theta}{2}$.

The total maximum angular error (accounting for the offset of $\Delta\theta/2$ in the initial state $|\psi\rangle$ from $|\beta\rangle$ in this case is,

$$\begin{aligned}\text{Total max. angular error} &= \frac{\pi}{4} + \text{Max. Initial Error} \\ &= \frac{\pi}{4} + \frac{\pi}{4} \frac{|\Delta\theta|}{\theta} \\ &= \frac{\pi}{4} \left(1 + \frac{|\Delta\theta|}{\theta} \right)\end{aligned}$$

ε_0 ,

Choosing $m = \left\lceil \frac{n}{2} \right\rceil + 1$,

$$|\Delta 0| \leq 2^{-m} = \frac{1}{2^m}$$

$$m = \left\lceil \frac{n}{2} \right\rceil + 1 \geq \frac{n}{2} + 1 \implies 2^m \geq 2^{\frac{n}{2} + 1}$$

$$\implies |\Delta 0| \leq 2^{-m} \leq 2^{-\frac{n}{2} - 1} = \frac{1}{2 \cdot 2^{\frac{n}{2}}}$$

Lemma :

$\sin x \leq x$ when $0 \leq x \leq \frac{\pi}{2}$

Proof

$$f(x) = \sin x - x, \quad 0 \leq x \leq \frac{\pi}{2}$$

$$f'(x) = \cos x - 1 \implies f'(0) = 1 - 1 = 0 \text{ \& } f'(\frac{\pi}{2}) = -1$$

$$f''(x) = -\sin x < 0 \quad \left. \vphantom{f''(x)} \right\} f(x) \text{ is a decreasing function in } [0, \frac{\pi}{2}]$$

$$f(0) = 0$$

$$f(\frac{\pi}{2}) = 1 - \frac{\pi}{2} < 0$$

$$\left. \vphantom{f(0)} \right\} \begin{aligned} f(x) &= \sin x - x \leq 0 \\ \therefore \sin x &\leq x \end{aligned}$$

Lemma $\Rightarrow \theta/2 \geq \sin \theta/2$ since $0 \leq \theta/2 \leq \pi/4$

$$\Rightarrow \frac{1}{\theta} \leq \frac{1}{\sin \theta}$$

$$\theta/2 \geq \sin \theta/2 = \sqrt{\frac{M}{2N}} \geq \sqrt{\frac{1}{2N}} \Rightarrow \theta/2 \geq \frac{1}{\sqrt{2N}} = \frac{1}{\sqrt{2} \times 2^{n/2}}$$

$$\frac{|\Delta \theta|}{\theta} = \frac{|\Delta \theta|}{2 \times \theta/2} \leq \frac{\frac{1}{2 \times 2^{n/2}}}{\frac{1}{\sqrt{2} \times 2^{n/2}}} = \frac{1}{\sqrt{2}} \leq \frac{1}{2}$$

If $M \leq N/2$, then $\sin \theta/2 = \sqrt{\frac{M}{N}} \leq \sqrt{\frac{1}{2}}$

$$\Rightarrow \theta/2 \geq \frac{1}{\sqrt{N}} = \frac{1}{2^{n/2}}$$

$$\Rightarrow \frac{|\Delta \theta|}{\theta} = \frac{|\Delta \theta|}{2 \times \theta/2} \leq \frac{\frac{1}{2 \times 2^{n/2}}}{\frac{1}{2^{n/2}}} = \frac{1}{2}$$

$$\therefore \frac{\pi}{4} \left(1 + \frac{|\Delta \theta|}{\theta}\right) = \frac{\pi}{4} \left(1 + \frac{1}{2}\right) = \frac{\pi}{4} \times \frac{3}{2} = \frac{3\pi}{8}$$

which is the max. angular error.

$$P(\text{measure } |f\rangle) \geq \sin^2\left(\frac{\pi}{2} - \theta/2\right) = \cos^2 \theta/2$$

where, $\theta/2$ is the angle b/w $G^k|\psi\rangle$ and $|f\rangle$.

The corresp. success probability is at least

$\cos^2(3\pi/8) \approx 0.15$ for the search algorithm.

The probability of obtaining an estimate of α this accurate is $5/6$.

$\therefore$ The total probability of obtaining a solution to the search problem is,

$$\frac{5}{6} \times \cos^2(3\pi/8) \approx 0.12$$

a probability which may quickly be boosted close to 1 by a few repetitions of the combined counting-search procedure.

$$P(\text{success} | \alpha) \geq \cos^2(3\pi/8) \approx 0.12$$

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